

45. If $Z(N, V, T)$ is the partition function of a system with N indistinguishable particles and volume V at temperature T and $Z(1, V, T)$ is the partition function of a system containing one particle in volume V at temperature T . Then, $Z(N, V, T)$ is

- (A) $[Z(1, V, T)]^N / N!$
 (B) $N[Z(1, V, T)] / N!$
 (C) $[Z(1, V, T)]^N$
 (D) $N[Z(1, V, T)]$

46. The probability of finding a system in a state with N_r particles and energy E_s is

$$P_{r,s} = \exp(-\alpha N_r - \beta E_s) / Z(z, V, T)$$

If z and Z are fugacity and grand partition function, $\beta = 1/kT$ and $\alpha = -\mu/kT$, the probability that the system in grand canonical ensemble has exactly N particles is

- (A) $Z(N, V, T) / Z(z, V, T)$
 (B) $z^N Z(N, V, T) / Z(z, V, T)$
 (C) $z^N Z(N, V, T) / [N Z(z, V, T)]$
 (D) $N z^N Z(N, V, T) / Z(z, V, T)$

47. A system has a fixed number of particles N . Suppose the energy scale is shifted by an arbitrary constant η so that the single particle levels are shifted by η . The new and the old partition functions, \tilde{Z} and Z will be related as,

- (A) $\tilde{Z} = Z$
 (B) $\tilde{Z} = e^{-\beta N / \eta} Z$
 (C) $\tilde{Z} = e^{-\beta \eta N} Z$
 (D) $\tilde{Z} = \eta e^{-\beta N} Z$

48. If S represents entropy of a macrostate and Ω , the number of microstates in it, then

- (A) S must be $> k \ln \Omega$
 (B) S must be $> k / \ln \Omega$
 (C) $S = k / \ln \Omega$
 (D) $S = k \ln \Omega$

k is Boltzmann constant.